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CERTAIN FINITE EXPANSIONS ASSOCIATED WITH A BASIC
ANALOGUE OF THE G-FUNCTION

ABSTRACT

In this paper the authors derive certain finite expansion formulae associated with a basic analogue of Meijer's G-function. The results proved in this paper are the extensions of the results given earlier by R. P. Agarwal [1] and N. Agarwal [2].

RESUMEN

En este trabajo los autores derivan ciertos desarrollos finitos asociados con una análoga básica de la función G de Meijer. Los resultados probados en este trabajo son las extensiones de los resultados dados anteriormente por R.P. Agarwal [1] y N. Agarwal [2].

INTRODUCTION AND PRELIMINARIES

Recently, Saxena, Modi and Kalla [12] defined the basic analogue of the G-function in the following manner:

$$G_{A,B}^{m_1, n_1} \left[z; q \begin{matrix} a_1, \dots, a_A \\ b_1, \dots, b_B \end{matrix} \right] =$$

$$\frac{1}{2\pi i} \int_C \frac{\prod_{j=1}^{m_1} G(q^{b_j - \xi}) \prod_{j=1}^{n_1} G(q^{1-a_j + \xi}) \prod z^\xi}{\prod_{j=m_1+1}^B G(q^{1-b_j + \xi}) \prod_{j=n_1+1}^A G(q^{a_j - \xi}) G(q^{1-\xi}) \sin \pi \xi} d\xi \quad (1)$$

where $0 \leq m_1 \leq B$, and $0 \leq n_1 \leq A$,

$$G(q^a) = \left\{ \prod_{n=0}^{\infty} (1 - q^{a+n}) \right\}^{-1}; \quad (2)$$

The conventional notations (a) and (a)+m will be employed to represent the following sequences of A parameters as mentioned below:

$$(a) = a_1, \dots, a_A; (a) + m = a_1 + m, \dots, a_A + m$$

Here the contour 'C' is a line parallel to $R_1(\omega\xi) = 0$ with indentations, if necessary, in such a manner, that all the poles of $G(q^{b_j - \xi})$ for $j=1, \dots, m_1$ are to the right and those of $G(q^{1-a_j + \xi})$ for $j=1, \dots, n_1$ to the left of it.

The integral converges, if

$$R_1 \{ \xi \log(z) - \log \sin \pi \xi \} < 0$$

for large values of $|\xi|$ on the contour, that is, if

$$| \left\{ \arg(z) - W_2 W_1^{-1} \log |z| \right\} | < \pi, \text{ where } |q| < 1, \log q = -W = -(W_1 + iW_2), W, W_1, W_2 \text{ are definite quantities, } W_1 \text{ and } W_2 \text{ being real.}$$

$$\text{Let } (a; q)_\mu = \prod_{j=0}^{\infty} \frac{(1-aq^j)}{(1-aq^{\mu+j})} \text{ for arbitrary } \mu \text{ and } a,$$

so that

$$(a; q)_n = \begin{cases} 1, & \text{for } n = 0 \\ (1-a)(1-aq) \dots (1-aq^{n-1}); & \forall n \in \{1, 2, 3, \dots\} \end{cases} \quad (3)$$

and

and

$${}_2\phi_1 \left[a, q^{-n}; c; q, cq^n/a \right] = \frac{(c/a; q)_n}{(c; q)_n} \quad (17)$$

the well poised ${}_6\phi_5$, Slater [13]

$${}_6\phi_5 \left[\begin{matrix} a, q\sqrt{a}, -q\sqrt{a}, b, c, q^{-n} \\ \sqrt{a}, -\sqrt{a}, aq/b, aq/c, aq^{1+n} \end{matrix}; q, aq^{1+n}/bc \right] = {}_r\phi_s \left[\begin{matrix} b_1^{-n+m}, b_2^{-n}, b_3^{-m}, \dots, b_r^{-m}; \\ a_1^{-m}, \dots, a_s^{-m} \end{matrix}; q, y \right] \quad (18)$$

$$= \frac{(aq; q)_n (aq/bc; q)_n}{(aq/b; q)_n (aq/c; q)_n} = {}_r\phi_s \left[\begin{matrix} b_1, \dots, b_r; \\ a_1, \dots, a_s \end{matrix}; q, y \right] \quad (21)$$

and the q-analogue of Dixon's theorem, namely,

$${}_4\phi_3 \left[\begin{matrix} a, -q\sqrt{a}, b, q^{-n} \\ -\sqrt{a}, aq/b, aq^{1+n} \end{matrix}; q, \sqrt{a} q^{1+n}/b \right] = \frac{(aq; q)_n (\sqrt{a} q/b; q)_n}{(\sqrt{a} q; q)_n (aq/b; q)_n} \quad (19)$$

where $y = x(1-q)^{s+1-r}$, which follows from the integral given by Watson [14], we obtain the following results.

$$\sum_{m=0}^n \frac{(q^{-n}; q)_m (b_1^{-n}, b_2^{-n}; q)_m (b_3; q)_m \dots (b_r; q)_m (yq^{2^{-n}})^m}{(q; q)_m (a_1; q)_m \dots (a_s; q)_m}$$

where $y = x(1-q)^{s+1-r}$

$$\sum_{m=0}^n \frac{(q^{-n}; q)_m (b_1; q)_m (b_3; q)_m \dots (b_r; q)_m (yq^{2^{-n}})^m}{(q; q)_m (a_1; q)_m \dots (a_s; q)_m}$$

SPECIAL CASES

(i) If we set $B = m_1 = r; n_1 = 0; A = s$ and replace

z by $\frac{-1}{x(1-q)^{s+1-r}}$ in (10) through (14), and apply the

relation

$$G_{s,r}^{r,0} \left[\begin{matrix} -1/y; q \\ (a) \\ (b) \end{matrix} \right] = \Pi \left[\begin{matrix} q, q^{a_1}, \dots, q^{a_s}; \\ b_1, \dots, b_r \end{matrix}; \right] = {}_r\phi_s \left[\begin{matrix} b_1, \dots, b_r; \\ a_1, \dots, a_s \end{matrix}; q, y \right] \quad (22)$$

$${}_r\phi_s \left[\begin{matrix} b_1, \dots, b_r; \\ a_1, \dots, a_s \end{matrix}; q, y \right] \quad (20)$$

where $y = x(1-q)^{s+1-r}$

$$\sum_{m=0}^n \frac{(q^{-n}; q)_m (b_2; q)_m \dots (b_r; q)_m (yq^n)^m}{(q; q)_m (a_1; q)_m \dots (a_s; q)_m}$$

$$= {}_{r+s} \phi_s \left[\begin{matrix} b_1+n, b_2+m, \dots, b_r+m; \\ a_1+m, \dots, a_s+m; \end{matrix} \middle| q, y \right]$$

$$= {}_{r+s} \phi_s \left[\begin{matrix} b_1, \dots, b_r; \\ a_1, \dots, a_s; \end{matrix} \middle| q, y/q^n \right] \quad (23)$$

where $y = x(1-q)^{s+1-r}$

$$\sum_{m=0}^n \frac{(q^{-n}; q)_m (a; q)_m (1 + \frac{a}{2}; q)_m (-1 - \frac{a}{2}; q)_m (b; q)_m}{(q; q)_m (a/2; q)_m (1+a-b+n; q)_m (-a/2; q)_m}$$

$$\frac{(b_1; q)_m \dots (b_r; q)_m (-y)^m q^{\frac{m}{2}(m+1)(a-b+n)m}}{(a_1; q)_m \dots (a_s; q)_m (1+a; q)_{2m}}$$

$$= {}_{r+2} \phi_{s+2} \left[\begin{matrix} 1+a-b+m, 1+a+n+m, b_1+m, \dots, b_r+m \\ a_1+m, \dots, a_s+m, 1+a-b+n+m, 1+a+2m; \end{matrix} \middle| q, y \right]$$

$$= {}_{r+s} \phi_s \left[\begin{matrix} b_1, \dots, b_r; \\ a_1, \dots, a_s; \end{matrix} \middle| q, y \right] \quad (24)$$

where $y = x(1-q)^{s+1-r}$

$$\sum_{m=0}^n \frac{(q^{-n}; q)_m (a; q)_m (-1 - \frac{a}{2}; q)_m (b_1; q)_m \dots (b_r; q)_m}{(q; q)_m (-a/2; q)_m (a_1; q)_m \dots (a_s; q)_m (1+a; q)_{2m}} (-y)^m q^{\frac{m}{2}(m+1)(a-b+n)m}$$

$$= {}_{r+1} \phi_{s+1} \left[\begin{matrix} 1+a+n+m, b_1+m, \dots, b_r+m; \\ a_1+m, \dots, a_s+m, 1+a+2m; \end{matrix} \middle| q, y \right]$$

$$= {}_{r+1} \phi_{s+1} \left[\begin{matrix} 1 + \frac{a}{2} + n, b_1, \dots, b_r; \\ a_1, \dots, a_s, 1 + \frac{a}{2}; \end{matrix} \middle| q, y \right] \quad (25)$$

where $y = x(1-q)^{s+1-r}$

(ii) On the other hand if we take $m_1 = B = r; n_1 = 0$ and $A = s$ in (10) and (13) and use the identity, Saxena, et al [12]

$$G_{s,r}^{r,0} \left[\begin{matrix} (a) \\ (b) \end{matrix} \middle| z; q \right] = E_q(r; b_j; s; a_i; z) \quad (26)$$

We arrive at the results given earlier by N. Agarwal, [1].

(iii) Next, if we take $m_1 = B = 2$ and $n_1 = 0 = A$ in (10), it gives a result due to R.P. Agarwal, [2]. If we further take $n=1$; (10) then reduces to another result given by Agarwal [2].

(iv) Finally, if we take $n = 1; m_1 = B = r; n_1 = 0$ and $A = s$ in (10) and (11), the known recurrence relations proved by Agarwal [1] are obtained.

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$$\text{where } L_n^{(\alpha)}(x|q) = \frac{[q^{1+\alpha}]_n}{[q]_n} {}_1\phi_1 \left[\begin{matrix} q^{-n}, q, xq^{1+\alpha+n} \\ q^{1+\alpha} \end{matrix} \right] \quad (1.6)$$

are the q-Laguerre polynomials of Hahn [1] satisfying the orthogonality condition (see Moak [2])

$$\int_0^\infty L_m^{(\alpha)}(x|q) L_n^{(\alpha)}(x|q) \frac{x^\alpha}{[-x]_\infty} dx = \frac{\Gamma(1+\alpha)\Gamma(-\alpha)[q^{1+\alpha}]_n}{\Gamma_q(-\alpha)[q]_n} q^{-n} (1-q)^{1+\alpha} \delta_n^m \quad (1.7)$$

where $\alpha > -1$ and δ_n^m is the familiar Kronecker delta.

2. ORTHOGONALITY

Using (1.7), in virtue of (1.4) and (1.5), one readily obtains the following orthogonality condition for the polynomials $H_{n,q}(x)$:

$$\int_0^\infty H_{n,q}(x) \cdot H_{m,q}(x) \frac{dx}{(-x^2q^2; q^2)_\infty} = \pi \sqrt{1-q^2} [q]_n q^{-(1+n)^2} \delta_n^m / \Gamma_{q^2}(1/2).$$

3. SOME PROPERTIES

From the explicit representation (1.2), we can obtain by usual series techniques; in view of q-binomial theorem [3, p.348, eq.(274)] and the elementary identity

$$[aq^{-n}]_\infty = (-1)^n a^n q^{-n(n+1)/2} [q/a]_n [a]_\infty,$$

the following generating function:

$$\sum_{n=0}^\infty \frac{[c]_n}{[q]_n} H_{n,q}(x) t^n = \frac{[cxt]_\infty}{[xt]_\infty} {}_3\phi_2 \left[\begin{matrix} 0, c, cq; q^2, -1/x^2 \\ q/xt, q^2/xt \end{matrix} \right] \quad (3.1)$$

In (3.1) replacing t by t/c and taking limit as $c \rightarrow \infty$ and changing again t by $-tq$, we get the generating function

$$\sum_{n=0}^\infty q^{n(n+1)/2} H_{n,q}(x) \frac{t^n}{[q]_n} = [-xtq]_\infty / [-t^2; q^2]_\infty \quad (3.2)$$

In (3.2) replacing x by xy and in the right hand side of resulting identity expanding the term $(-y^2t^2; q^2)_\infty / (-t^2; q^2)_\infty$ by q-binomial theorem and then equating the coefficients of t^n on both sides, we get a multiplication formula

$$H_{n,q}(xy) = \sum_{j=0}^{[n/2]} (-1)^j \frac{[q^{-n}]_{2j}}{(q^2; q^2)_j} y^{n-2j} (y^2; q^2)_j H_{n-2j,q}(x).$$

$$\text{Since } [x]_\infty = \sum_{n=0}^\infty (-1)^n q^{n(n-1)/2} \frac{x^n}{[q]_n},$$

from (3.2), by routine method, we have expansion of x^n in the form

$$x^n = \sum_{r=0}^{[n/2]} \frac{[q]_n}{(q^2; q^2)_r [q]_{n-2r}} q^{r(3r-2n-2)} H_{n-2r,q}(x).$$

Again, from the generating function (3.2)

$$\delta \sum_{n=0}^\infty q^{n(n+1)/2} H_{n,q}(x) \frac{t^n}{[q]_n} = tq \frac{[-xtq^2]_\infty}{(-t^2; q^2)_\infty} \quad (3.3)$$

so that

$$\sum_{n=0}^\infty q^{n(n+1)/2} \delta H_{n,q}(x) \frac{t^n}{[q]_n} = tq \sum_{n=0}^\infty q^{n(n+1)/2} H_{n,q}(xq) \frac{t^n}{[q]_n},$$

which yields $\delta H_{0,q}(x) = 0$, and for $n \geq 1$,

$$\delta H_{n,q}(x) = (1-q^n) q^{1-n} H_{n-1,q}(xq) \quad (3.4)$$

or, more generally,

$$\delta^m H_{n,q}(x) = [q^{n-m-1}]_m q^{m(m-n)} H_{n-m,q}(xq^m), \quad n \geq m \geq 0.$$

Alternatively, we may write (3.3) as

$$\delta \sum_{n=0}^{\infty} q^{n(n+1)/2} H_{n,q}(x)(1+t^2) \frac{t^n}{[q]_n}$$

$$= tq \sum_{n=0}^{\infty} q^{n(n+1)/2} H_{n,q}(x) \frac{t^n q^n}{[q]_n}$$

which readily yields,

$$\delta H_{n,q}(x) + q^{1-2n} (1-q^{n-1}) (1-q^n) \delta H_{n-2,q}(x) = (1-q^n) H_{n-1,q}(x).$$

It is easy to verify that

$$\left(x - \frac{t}{q}\right) \delta \left\{ \frac{[-xtq]_{\infty}}{(-t^2; q^2)_{\infty}} \right\} - t \delta_t \left\{ \frac{[-xtq]_{\infty}}{(-t^2; q^2)_{\infty}} \right\} = 0,$$

which, in conjunction with (3.2), gives

$$xq^{1+n} \delta H_{n,q}(x) = q^{1+n} (1-q^n) H_{n,q}(x) + (1-q^n) \delta H_{n-1,q}(x) \quad (3.5)$$

Using (3.4) in (3.5), we have

$$q^n H_{n,q}(x) - xq H_{n-1,q}(xq) + q^{1-n} (1-q^{n-1}) H_{n-2,q}(xq) = 0. \quad (3.6)$$

From (3.4), in view of definition (1.1), we get

$$q^n H_{n,q}(x) = q^n H_{n,q}(xq) + xq(1-q^n) H_{n-1,q}(xq). \quad (3.7)$$

Notice, however, that by eliminating the term $q^n H_{n,q}(x)$ between (3.6) and (3.7), and then changing x to x/q , we at once obtain the pure recurrence relation

$$H_{n,q}(x) - xH_{n-1,q}(x) + q^{1-2n} (1-q^{n-1}) H_{n-2,q}(x) = 0.$$

Further, in (3.5) replacing x by xq and appealing to (3.4), we get a q -difference equation for q -Hermite polynomials in the form:

$$\delta^2 H_{n,q}(x) - xq^3 \delta H_{n,q}(xq) + q^3 (1-q^n) H_{n,q}(xq) = 0. \quad (3.8)$$

Following Moak [2,p.34], on account of (3.8), it is not hard to establish

$$\delta^2 u(x) + \left\{ \frac{1+q}{qx^2} - \frac{(1+qx^2q^{4+n})(-x^2q^6; q^4)_{\infty}}{qx^2(-x^2q^4; q^4)_{\infty}} \right\} u(xq) = 0,$$

where $u(x) = H_{n,q}(x)/(-x^2q^4; q^4)_{\infty}$

Letting $b \rightarrow 0$ in the q -extension of Euler's transformation [3,p.348, eq.(281)], we obtain

$${}_2\phi_1 \left[\begin{matrix} 0, a; q, z \\ c \end{matrix} \right] = \frac{1}{[z]_{\infty}} {}_1\phi_1 \left[\begin{matrix} c/a; q, -az \\ c \end{matrix} \right], \quad (3.9)$$

which is a q -extension of Kummer's first formula [3,p.322,eq.(183)].

Since $\frac{1}{[x]_{\infty}} = \sum_{n=0}^{\infty} \frac{x^n}{[q]_n}$, we have

$$\delta^{2n} \left\{ \frac{1}{(-x^2q^2; q^2)_{\infty}} \right\} = \sum_{r=n}^{\infty} \frac{(-1)^r q^{2r}}{(q^2; q^2)_r} x^{2r-2n} [q^{2r-2n+1}]_{2n},$$

which on replacing r by $n+r$ and doing some straightforward manipulations reduce to

$$\delta^{2n} \left\{ \frac{1}{(-x^2q^2; q^2)_{\infty}} \right\} = (-1)^n q^{2n} (q; q^2)_n {}_2\phi_1 \left[\begin{matrix} 0, q^{1+2n}; q^2, -x^2q^2 \\ q \end{matrix} \right] \quad (3.10)$$

Now, by applying (3.9) to (3.10) and then using (1.4) and (1.6), we find that

$$H_{2n,q}(x) = (-1)^{2n} q^{-3n-2n^2} (-x^2q^2; q^2)_{\infty} \delta^{2n} \left\{ \frac{1}{(-x^2q^2; q^2)_{\infty}} \right\} \quad (3.11)$$

Similarly, we can obtain

$$H_{2n+1,q}(x) = (-1)^{2n+1} q^{-2-5n-2n^2} (-x^2 q^2; q^2)_\infty \delta^{2n+1} \left\{ \frac{1}{(-x^2 q^2; q^2)_\infty} \right\} \quad (3.12)$$

Combining (3.11) and (3.12), we get a Rodrigues type representation for the polynomials $H_{n,q}(x)$ as

$$H_{n,q}(x) = (-1)^n q^{-n(n+3)/2} (-x^2 q^2; q^2)_\infty \delta^n \left\{ \frac{1}{(-x^2 q^2; q^2)_\infty} \right\}$$

More generally, one can obtain

$$\delta^k \left\{ \frac{H_{n,q}(x)}{(-x^2 q^2; q^2)_\infty} \right\} = (-1)^k q^{nk + (1/2)k(3+k)} \frac{H_{n+k,q}(x)}{(-x^2 q^2; q^2)_\infty}$$

which, for $k=1$, reduces to a recurrence relation

$$xq^{n+2} H_{n+1,q}(x) + H_{n,q}(x) - (1+x^2 q^2) H_{n,q}(xq) = 0.$$

One notes that all the properties for the polynomials $H_{n,q}(x)$ reduce, in view of (1.3), to corresponding properties for the Hermite polynomials $H_n(x)$.

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